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The proposed solution is seen directly to $\frac{dx}{\lambda} = \frac{dy}{\lambda} = \frac{dz}{\lambda}$ satisfy this and to satisfy the boundary

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conditions $B \left| \pm \delta \right| = B_a$. (b) For $\delta \ll \lambda L$, $\left(\frac{2}{\lambda} \right) \frac{1}{2} \left(\frac{x}{\lambda} \right) \cosh = 1 + \left| \frac{x}{\lambda} \right| + \dots \frac{2}{\lambda L} \left(\frac{x}{\lambda} \right) \frac{1}{2} \left(\frac{\delta}{\lambda} \right) \cosh = 1 + \left| \frac{x}{\lambda} \right| + \dots \frac{2}{\lambda L} \left(\frac{x}{\lambda} \right) \left(\frac{\delta}{\lambda} \right)$ therefore $B(x) = B_a - B_a \frac{1}{8} \frac{x^2}{\lambda^2} - 4x^2 \cdot \frac{2}{\lambda^2}$.

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$S(\text{basis}) = 1 + e^{-i\pi} = 0$. Now $S(\text{fcc}) = 0$ only if all indices are even or all indices are odd. If all indices are even the structure factor of the basis vanishes unless $v_1 + v_2 + v_3 = 4n$, where n is an integer. For example, for the reflection (222) we have $S(\text{basis}) = 1 + e^{-i3\pi} = 0$, and this reflection is forbidden.

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The projection of the magnetic moment of the upper state along the field direction is $-J-t$ and of the lower state is $J-t$. The resultant magnetization for N atoms per unit volume is M , with $x = J-tB/kBT$, $e^x - e^{-x}$. $M = (N\uparrow - N\downarrow)J-t = N\uparrow J-t - N\downarrow J-t = N\uparrow J-t \frac{e^x - e^{-x}}{e^x + e^{-x}}$. (17) For $x \sim 1$, $\tanh x = x$, and we have $M = N\uparrow J-t(J-tB/kBT)$ (18)

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cube. If θ is the angle between them, their scalar product gives $\cos \theta = -1/3$, whence. $\theta = \cos^{-1} 1/3 = 90^\circ + 19^\circ 28' = 109^\circ 28'$. 2. The plane (100) is normal to the x axis. It intercepts the a' axis at $2a'$ and the c' axis. at $2c'$; therefore the indices referred to the primitive axes are (101). Similarly, the plane.

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